Implicit eddy parameterization for Quasi-geostrophic models



Louis Thiry, ERC STUOD Postdoc, IRMAR, INRIA, ODISSEY Team

Joint work with Long Li, Etienne Mémin (INRIA) and Guillaume Roullet (UBO, LOPS).

Multi-layer quasi-geostrophic model

- *n* stacked layer, thickness H_k and density ρ_k .
- State variables:

pressure: $\mathbf{p} = (p_1(x, y), \dots, p_n(x, y))$ potential vorticity: $\mathbf{q} = (q_1(x, y), \dots, q_n(x, y))$

$$\partial_t \mathbf{q} + (\mathbf{u} \cdot \nabla) \mathbf{q} = 0$$

$$\Delta_H \mathbf{p} - f_0^2 A \mathbf{p} = f_0 \mathbf{q} - f_0 \beta y$$

$$- f_0 \mathbf{u} = \partial_y \mathbf{p}, \qquad f_0 \mathbf{v} = \partial_x \mathbf{p}$$

$$A = \begin{bmatrix} \frac{1}{H_1 g'_1} & \frac{-1}{H_1 g'_1} & \cdot & \cdot \\ \frac{-1}{H_2 g'_1} & \frac{1}{H_2} \left(\frac{1}{g'_1} + \frac{1}{g'_2} \right) & \frac{-1}{H_2 g'_2} & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \frac{-1}{H_n g'_{n-1}} & \frac{1}{H_n g'_{n-1}} \end{bmatrix} .$$

Idealized double-gyre configuration

- 3 layers, $H_k = 350, 750, 2900$ m.
- Rectangular domain, 3480×4800 km, solid boundaries, no-slip b.c.
- Baroclinic Rossby radii: 39, 22 km.
- Linear bottom drag, idealized wind stress on top, magnitude $\tau_0 = 2e^{-5}m^2s^{-2}$



< 注 > …

Numerical implementation

Following Hogg et al. (2014)

- Usual 5-points laplacian discretization.
- Advection with 9-points energy-conserving Arakawa Jacobian.
- Additional hyperviscosity set with Munk scale.
- Elliptic equation solved with type-I Discrete Sine Transform.
- Heun-RK2 time stepping.

$$\partial_t \mathbf{q} = \frac{1}{f_0} J(\mathbf{q}, \mathbf{p}) - \frac{\partial_4}{f_0} \Delta^3 \mathbf{p} + \text{forcing} + \text{drag}$$
$$(\Delta - f_0^2 A) \mathbf{p} = f_0 \mathbf{q} - f_0 \beta (y - y_0),$$

https://github.com/louity/qgm_pytorch

Eddy-resolving regime

- 769×961 grid, resolution 5km, dt=800s.
- Hyperviscosity $a_4 = 2 \ 10^9 \ m^4 s^{-1}$
- Apparition of proper eastward jet.
- Animation



Layers zonal velocity, 5km, 40y spinup

Eddy-resolving regime

• Rich meso-scale eddies field in the recirculation zone.



Layers relative vorticity, 5km, 40y spinup

Eddy-permitting regime

- 193×241 grid, resolution 20km, dt=1200s.
- Hyperviscosity $a_4 = 8 \ 10^{10} \ m^4 s^{-1}$
- No proper eastward jet.
- Animation



Layers zonal velocity, 20km, 40y spinup

Eddy-permitting regime

• Almost no eddies.



Layers relative vorticity, 20km, 40y spinup

Non eddy-resolving regime

- 97×121 grid, resolution 40km, dt=1400s.
- Hyperviscosity $a_4 = 5 \ 10^{11} \ m^4 s^{-1}$
- Tiny eastward jet without any eddy around.
- Animation



 \longrightarrow Need for eddy parameterizations.

Existing parameterizations

Eddy parameterization for QG models tested on double-gyre configuration:

- Zanna et al. (2017): deterministic + stochastic.
- Berloff et al. (2021): deterministic data-driven.
- Li et al. (2020): stochastic + mean term (= deterministic).
- Uchida et al. (2022): deterministic.

Deterministic parameterizations help producing large scale structures (jet).

Stochasticity improves variability, finer-scale structures and ensemble spread.

 \longrightarrow Importance of good deterministic parameterization as basis for stochastic ones.

Implicit parameterization

Roullet and Gaillard (2022): A fast monotone discretization of the rotating shallow water equations.

"Monotone? Because what is the point of invoking an adhoc dissipation or a sophisticated SGS theory when a good numerics can do both?" Analogous of implicit-LES for eddy parameterizations.

Implicit parameterization

Roullet and Gaillard (2022): A fast monotone discretization of the rotating shallow water equations.

"Monotone? Because what is the point of invoking an adhoc dissipation or a sophisticated SGS theory when a good numerics can do both?" Analogous of implicit-LES for eddy parameterizations.

Ingredients:

- **p**, **q** staggered grid
- Finite volume for PV and material conservation.
- High-order WENO (Balsara et al., 2016) interpolation for advection.
 ⇒ implicit diffusion replaces hyper-viscosity.
- Stable strongly preverving RK3 time-stepping.

Numerical implementation



Figure: Staggered grid discretization

・ 同 ト ・ ヨ ト ・ ヨ ト

Results in eddy-permitting regime

- 193×241 grid, resolution 20km, dt=2000s.
- No hyper-viscosity.
- Half-length eastward jet.
- Animation



Layers zonal velocity, 20km, 40y spinup

Results in eddy-permitting regime

• Large meso-scale eddies in the recirculation zone.

Top-layer relative vorticity, 20km, 40y spinup



■▶▲■▶ ■ ⊘⊘⊘

14 / 28

- 97×121 grid, resolution 40km, **dt=4000**s.
- No hyper-viscosity.
- Third-length eastward jet.
- Animation



15 / 28

- 769×961 grid, resolution 5km, dt=800s.
- Symmetry breaking and effective resolution

Standard QG Our method 0+0

Top-layer relative vorticity, 5km, 1.5y spinup

• Symmetry breaking and effective resolution

Top-layer relative vorticity, 5km, 1.5y spinup
Standard QG
Our method



• Symmetry breaking and effective resolution

Top-layer relative vorticity, jet-region, 5km, 40y spinup







Statistics



Arakawa (q,p) statistics

・ロト・西ト・西ト・西ト・日・ 今々ぐ

WENO5 implicit parameterization

- Works in the three regimes.
- Shows different advantages.
- Removes the viscosity CFL condition: larger dt.

WENO5 implicit parameterization

- Works in the three regimes.
- Shows different advantages.
- Removes the viscosity CFL condition: larger dt.
- ... but can still be improved.

q to p interpolation

$$(\Delta - f_0^2 A)\mathbf{p} = \operatorname{Interp}_{q \to p} (f_0 \mathbf{q} - f_0 \beta(y - y_0))$$

- Interpolation needed to solve elliptic equation
- 4-points interpolation has bad frequency response
 - \implies high-frequency are discarded before solving elliptic equation



Figure: Staggered grid discretization

Interpolation with DST-II

$$(\Delta - f_0^2 A)\mathbf{p} = \operatorname{Interp}_{q \to p} (f_0 \mathbf{q} - f_0 \beta(y - y_0))$$

- Elliptic equation solved with DST-I
- DST-II(\cdot) = DST-I(Spectral-Interp_{q $\rightarrow p$}(\cdot)) Spectral interpolation \implies highest possible order



Figure: DST-I and DST-II equivalence on square grid

Eddy-resolving regime:

- 769×961 grid, resolution 5km, dt=800s.
- Animation.

Eddy-resolving regime:

- 769×961 grid, resolution 5km, dt=800s.
- Animation.

• Longer eastward jet.

Eddy-permitting regime:

- 193×241 grid, resolution 20km, dt=2000s.
- Animation

Eddy-resolving regime:

- 769×961 grid, resolution 5km, dt=800s.
- Animation.

• Longer eastward jet.

Eddy-permitting regime:

- 193×241 grid, resolution 20km, dt=2000s.
- Animation
- Proper eastward jet.

Non eddy-resolving regime:

- 97×121 grid, resolution 40km, dt=4000s.
- Animation

Eddy-resolving regime:

- 769×961 grid, resolution 5km, dt=800s.
- Animation.

Longer eastward jet.

Eddy-permitting regime:

- 193×241 grid, resolution 20km, dt=2000s.
- Animation
- Proper eastward jet.

Non eddy-resolving regime:

- 97×121 grid, resolution 40km, dt=4000s.
- Animation

ALIASING!

Spectral methods need anti-aliasing treatment, e.g. Dedalus.

Conclusions

Our implicit eddy parameterization

- resolves large-scale structures (eddy-permitting and non-eddy-resolving regimes).
- increases effective resolution, accelerates spin up and symmetry breaking (eddy-resolving regime).
- works in a wide range of resolutions: 1.5 to 40km.
- allows removing ad-hoc dissipation, hence the viscosity CFL condition.
- is complementary to physical or stochastic parameterizations.

Take home message

• At low resolutions, preserving high-frequencies is crucial...

Conclusions

Our implicit eddy parameterization

- resolves large-scale structures (eddy-permitting and non-eddy-resolving regimes).
- increases effective resolution, accelerates spin up and symmetry breaking (eddy-resolving regime).
- works in a wide range of resolutions: 1.5 to 40km.
- allows removing ad-hoc dissipation, hence the viscosity CFL condition.
- is complementary to physical or stochastic parameterizations.

Take home message

- At low resolutions, preserving high-frequencies is crucial...
- ...but can be tricky : Aliasing ! Rossby waves dissipation ? Badly resolved western boundary-layer ?

Conclusions

Our implicit eddy parameterization

- resolves large-scale structures (eddy-permitting and non-eddy-resolving regimes).
- increases effective resolution, accelerates spin up and symmetry breaking (eddy-resolving regime).
- works in a wide range of resolutions: 1.5 to 40km.
- allows removing ad-hoc dissipation, hence the viscosity CFL condition.
- is complementary to physical or stochastic parameterizations.

Take home message

- At low resolutions, preserving high-frequencies is crucial...
- ...but can be tricky : Aliasing ! Rossby waves dissipation ? Badly resolved western boundary-layer ? (Suggestions ?)

Perspectives

- Add stochastic parameterization: Mémin (2014).
- Jump to WENO7 ? Or Semi-lagrangian ?
- Ability to resolve eddies:
 - Realistic wind
 - Inclusion of surface buoyancy: Lapeyre and Klein (2006) (animation)
- More realistic simulation :
 - Move to non-trivial geometries.
 - Beyond quasi-geostrophy: shallow-water ?
- Suitable model for data assimilation ?

Questions ?

★ E ► < E ►</p>

References I

- Dinshaw S Balsara, Sudip Garain, and Chi-Wang Shu. An efficient class of weno schemes with adaptive order. *Journal of Computational Physics*, 326:780–804, 2016.
- Pavel Berloff, Evgeny Ryzhov, and Igor Shevchenko. On dynamically unresolved oceanic mesoscale motions. *Journal of Fluid Mechanics*, 920, 2021.
- AM Hogg, JR Blundell, WK Dewar, and PD Killworth. Formulation and users' guide for q-gcm, 2014.
- GUILLAUME Lapeyre and PATRICE Klein. Dynamics of the upper oceanic layers in terms of surface quasigeostrophy theory. *Journal of physical oceanography*, 36(2):165–176, 2006.
- Long Li, Etienne Mémin, and Bertrand Chapron. Quasi-geostrophic flow under location uncertainty. In *Seminar of Stochastic Transport in Upper Ocean Dynamics (STUOD) project*, pages 1–52, 2020.

References II

- Etienne Mémin. Fluid flow dynamics under location uncertainty. Geophysical & Astrophysical Fluid Dynamics, 108(2):119–146, 2014.
- Guillaume Roullet and Tugdual Gaillard. A fast monotone discretization of the rotating shallow water equations. *Journal of Advances in Modeling Earth Systems*, 14(2):e2021MS002663, 2022.
- Takaya Uchida, Bruno Deremble, and Stephane Popinet. Deterministic model of the eddy dynamics for a midlatitude ocean model. *Journal of Physical Oceanography*, 2022.
- Laure Zanna, PierGianLuca Porta Mana, James Anstey, Tomos David, and Thomas Bolton. Scale-aware deterministic and stochastic parametrizations of eddy-mean flow interaction. *Ocean Modelling*, 111:66–80, 2017.

< ⊒ >